

Saddle Coil Design for MRI/MRS at 14.1 Tesla: A Favorable Alternative to the Quadrature Birdcage Coil

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Synopsis

A single channel ¹H saddle coil and a quadrature ¹H 8-leg birdcage coil were compared in terms of B_1^+ homogeneity and B_1^+ efficiency for preclinical small-animal imaging at 14.1 Tesla. Electromagnetic field simulations showed that the saddle coil generated a more homogeneous B_1^+ field for all regions of interest (up to 13.7 mm central disk diameter) than the birdcage coil. Saline phantom measurements showed a more constant flip angle excitation upon multi-slice imaging in axial direction and a 66% higher B_1^+ efficiency for the saddle coil.

Introduction

Many applications in MRI rely on high transmit field (B_1^+) homogeneity. Simultaneously, it is advantageous to maintain high B_1^+ efficiency and receive sensitivity. To insure B_1^+ homogeneity, volume coils are preferred to surface coils. For preclinical research, thus small regions of interest (ROI), Birdcage Coils (BC) and saddle coils are popular design choices¹. However, BC used at ultra-high field such as 14.1T have to deal with rather small diameters due to the reduced bore size, which entails distorted B_1^+ homogeneity^{2,3}. Previous work showed that a single channel saddle coil yielded higher sensitivity than a quadrature BC for an object of the size of a rat upper thigh⁴. In the present work, we investigated how saddle coil and BC perform at 14.1T in terms of B_1^+ homogeneity and B_1^+ efficiency. Quantitative analysis was done via electromagnetic field simulations and phantom measurements.

Methods

Both coils were designed to accommodate a rat head or thigh and to fit the 14.1T magnet bore, with coil diameter $\phi=50.6$ mm, length $l=27$ mm for the shielded 8-leg ¹H quadrature BC and $\phi=34$ mm, $l=22$ mm for the unshielded saddle coil (more compact due to its more open design)⁴. The saddle coil angular aperture $\alpha=120^\circ$ was adjusted to minimize radially dependent contributions to the central magnetic field⁵. The electromagnetic field simulations were performed using finite-difference time-domain model based software Sim4Life4.0.1 (ZMT, Switzerland).

Flip angle (θ) maps of a saline phantom were acquired via gradient echo imaging (Gaussian pulses of $\tau=2$ ms duration, $TR/TE=20000/4.83$ ms, $\theta=60^\circ/120^\circ$, matrix size = 128×128 , $FOV=25 \times 25$ mm²) and calculated according to the double angle method, $\theta = \cos^{-1}(I_2/I_1)^{0.7}$ with voxel-wise magnitude image intensity I . B_1^+ homogeneity was quantified over a ROI of 11mm diameter (ROI₀₁₁) using: (1) Non-Uniformity (NU) ($=\sigma(B_1^+)/B_1^+$) and (2) Relative Uniformity (RU) (percentile number of voxels which deviate by $\leq 10\%$ from B_1^+)⁶. The calibration power was determined for localized (11.2x11.2x3mm³ voxel) acquisition using asymmetric 90° stimulated echo acquisition mode pulses (width $\tau=500$ s, $TE=2.8$ ms)^{9,10}. Resonance modes and quality factors (Q) were measured on a network analyzer¹¹.

Results

Fig.1 displays scattering parameters and respective resonance modes. The saddle coil has a single resonance at 600MHz contrarily to the BC which has a number of modes. Transmit field simulations are shown in Fig.2. Homogeneity over axial slices spanning a 14mm region were non-uniformity $NU^*=1.0 \pm 0.1\%$ and $NU=2.7 \pm 0.1\%$ for saddle coil and BC, respectively. For ROI₀₁₁ we found $NU^*=1.7\%$, $RU^*=100\%$ for the saddle coil and $NU=4.0\%$, $RU=95.4\%$ for the BC on the axially central slice. The flip-angle dependency on the axial slice position showed a constant increase from $\alpha|_{z=-6\text{mm}}=20.4 \pm 1.0^\circ$ to $\alpha|_{z=+6\text{mm}}=33.6 \pm 1.1^\circ$ in an 11mm diameter central ROI (147 voxel) for the BC (Fig.3). The saddle coil created constant nominal flip angles over the first 4 slices in direction of the source with $\alpha|_{-6 < z < 0\text{mm}}=30.2 \pm 0.6^\circ$ and only then gradually decreased to $\alpha|_{z=+6\text{mm}}=19.0 \pm 2.4^\circ$. Fig.4 shows the transmit field in phantoms where $NU^*=5.0\%$, $RU^*=91.1\%$ for the saddle coil and $NU=3.5\%$, $RU=97.7\%$ for the BC. The saddle coil needed 5-6dB less power to create a 90° pulse upon localized acquisition. B_1^+ efficiency was estimated to be $8.8 \pm 0.6 \mu T / \sqrt{kW}$ for saddle coil and $5.3 \pm 0.2 \mu T / \sqrt{kW}$ for BC. Table1 contains unloaded (Q_{ul}) and loaded (Q_l) quality factors with $Q_{ul}/Q_l=2.7$ and $Q_{ul}/Q_l=1.1$ for saddle coil and BC, respectively.

Discussion

Electromagnetic field simulations showed the saddle coil to outperform the BC in terms of transmit field homogeneity over (a) slices in axial direction and (b) over ROIs with different diameters in the axially central slice. The same pattern of slight B_1^+ inhomogeneity derived from simulations was also observed in phantom measurements (diagonal under-intensity shades in figures 2A,B and 4A,B), supporting the validity of the simulations results. Inaccuracies between the simulations and measurements arise from imperfections in angular symmetry, non-ideal current distribution throughout the rungs, RF losses and balanced segmenting capacitors. On the practical side, tuning and matching is not trivial for BC due to coupling between the two channels, multi-modal resonance spectra and potential splitting of degenerate modes. It is all the more straightforward in the case of a single-resonance saddle coil. The ratio $Q_{ul}/Q_l=1.1$ for the BC was confirmed in previous publications¹². Due to the lower number of RF components on the saddle coil, less power is dissipated in its lumped components. This is reflected in a higher ratio $Q_{ul}/Q_l=2.7$ and consequently in the power requirements for a 90° pulse. The difference of 5-6dB means that the BC required 2x more power to generate a 90° flip angle. The saddle coil had a 66% higher coil efficiency than the BC.

Conclusion / Outlook

We conclude that the designed saddle coil provides a higher transmit field homogeneity than the BC coil and requires less input power to generate a given flip angle with a substantially higher B_1^+ efficiency. This provides the possibility to generate higher B_1^+ values with the same RF input power, enabling the use of shorter pulses with larger bandwidth and/or a better pulse response profile for a given flip angle and it opens the door to short-TE MRS. In this respect, this saddle coil geometry yields considerable advantages at 14.1T for combinations of advanced MRI approaches requiring strong B_1^+ homogeneity with spectroscopic studies or whenever power limitations are encountered.

Acknowledgements

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Figures

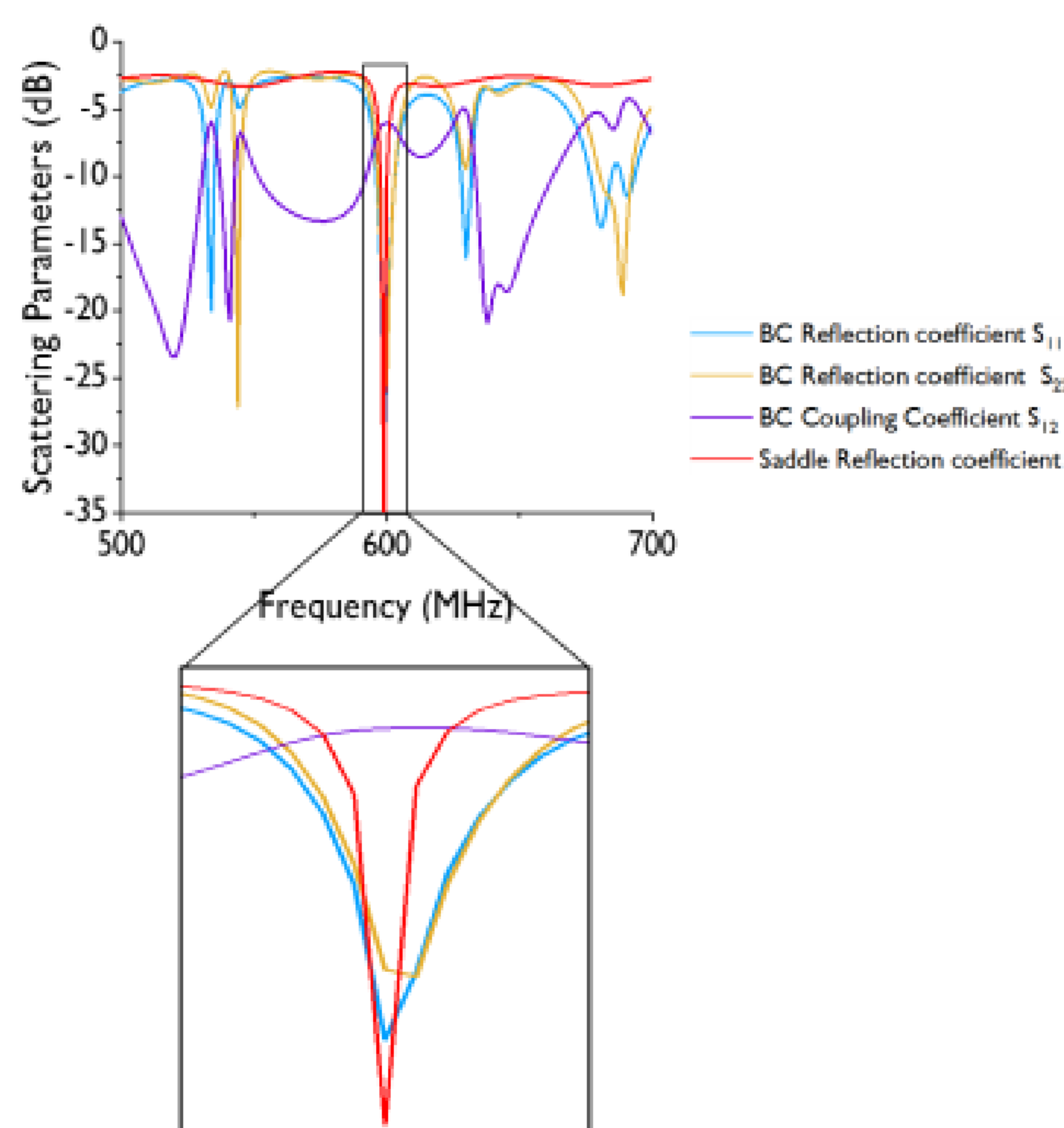


Figure 1: Scattering parameters for loaded coils within a 200 MHz window around the resonance frequency. Reflection and transmission coefficients (the latter is only relevant for the quadrature BC coil) show the different resonance modes.

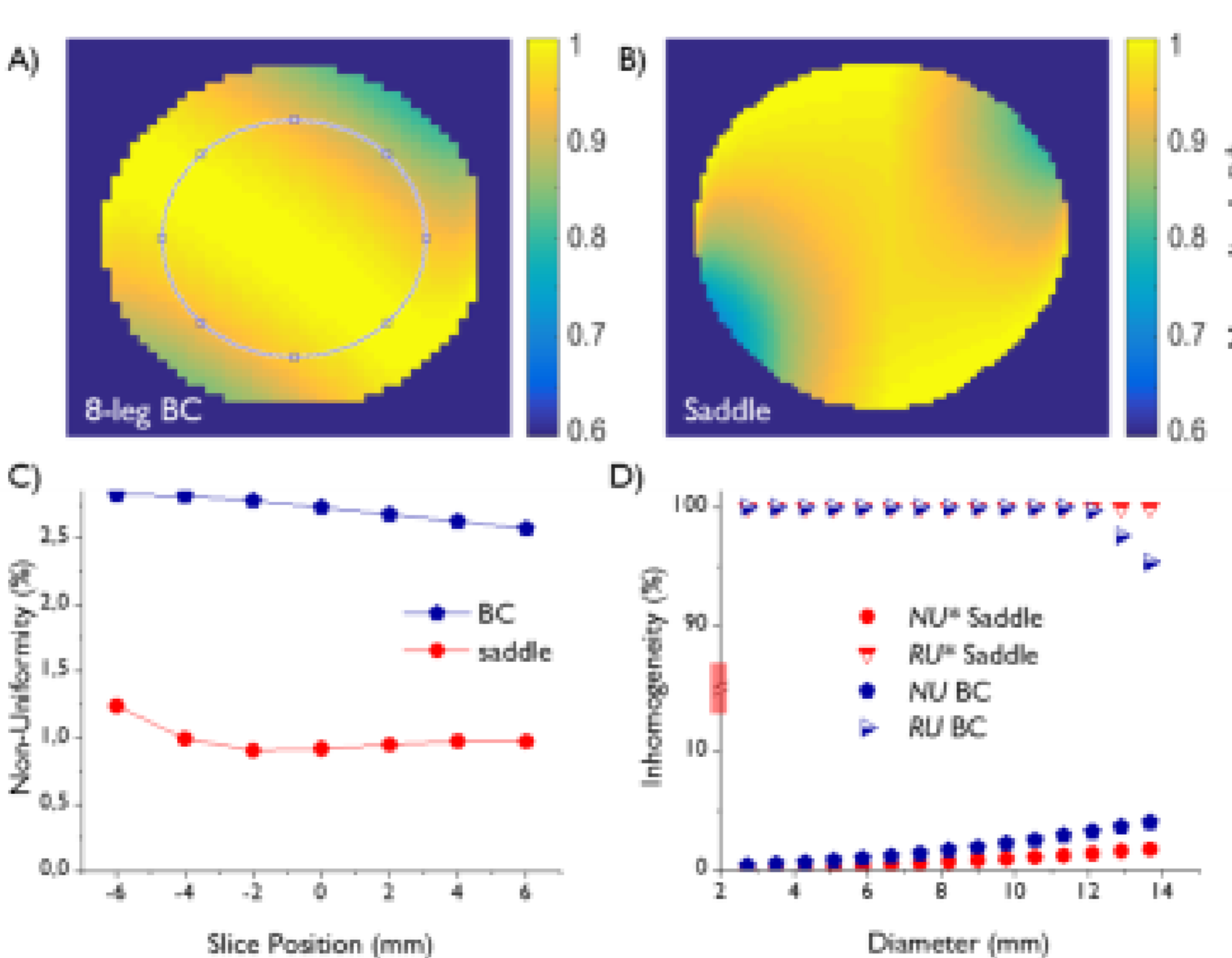


Figure 2: Simulated and normalized B_1^+ maps obtained from voxel-wise calculation $B_1^+ = \theta/\gamma\tau\epsilon$, with ¹H gyromagnetic ratio γ and normalized pulse envelope correction factor $\epsilon = 0.413$ for the applied Gaussian pulse. Axially centered B_1^+ maps generated by **A)** 8-leg BC and **B)** saddle coil. **C)** Dependence of NU on the slice distance from the isocenter. Statistics were calculated within an 11 mm diameter ROI (encircled in A). **D)** NU and RU of the B_1^+ maps as a function of the ROI diameter on the axially central slice (note the cut of the y-axis).

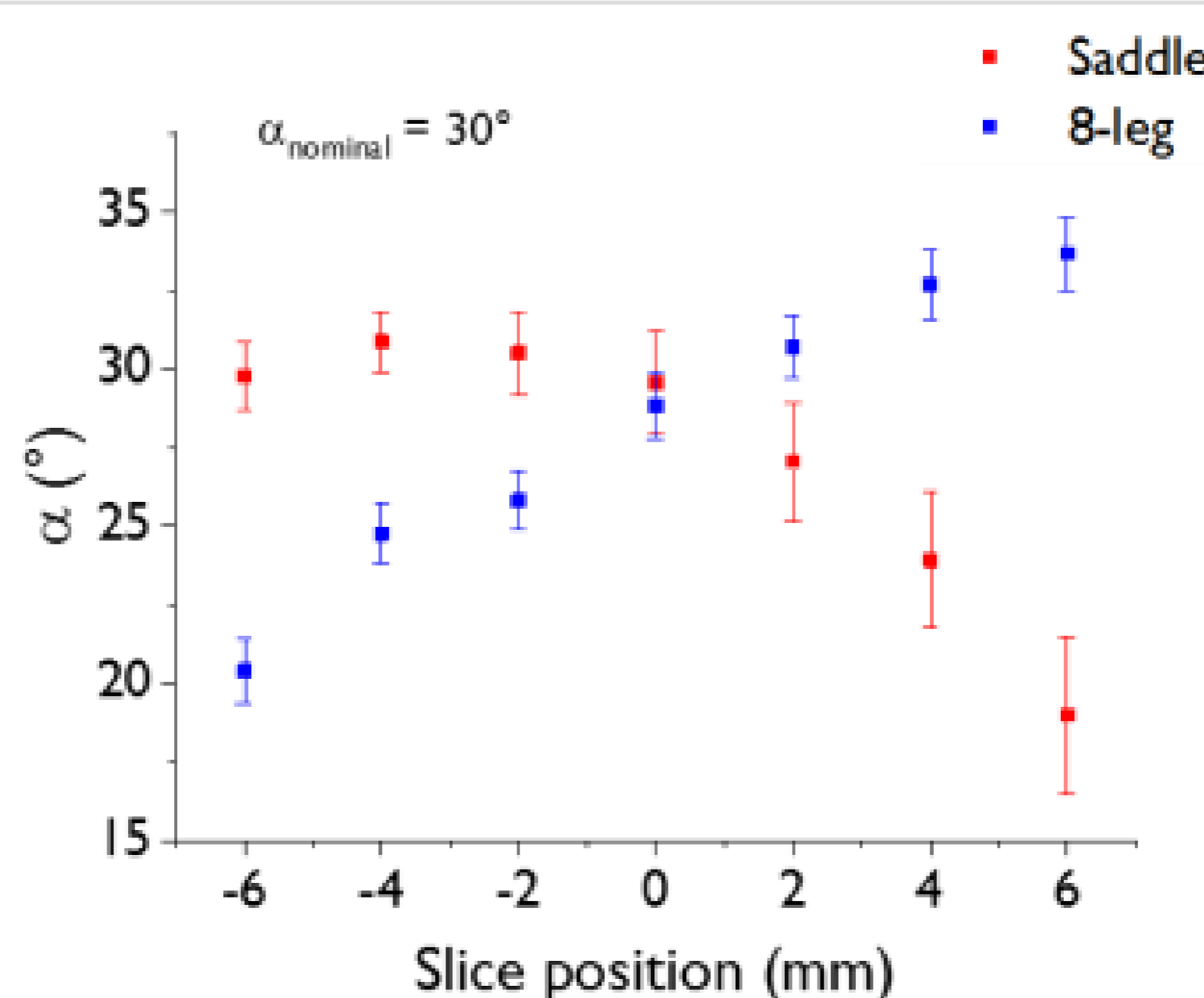


Figure 3: Mean effective flip angle α over the excited slice as a function of slice position in z-direction (B_1^+) for α calibrated to 30° in the center slice (position = 0 mm). Error bars denote the standard deviation of the voxel-wise flip angles over a 11 mm diameter central ROI (147 voxel). The same gradient echo sequence was used as for power calibration measurements: $TR = 10$ s, $TE = 2.6$ ms, slice thickness = 2 mm, $FOV = 25 \times 25$ mm², matrix size = 32×32 . The scale is chosen such that the origin is in the isocenter and positive slice positions are in direction of the coil current source.

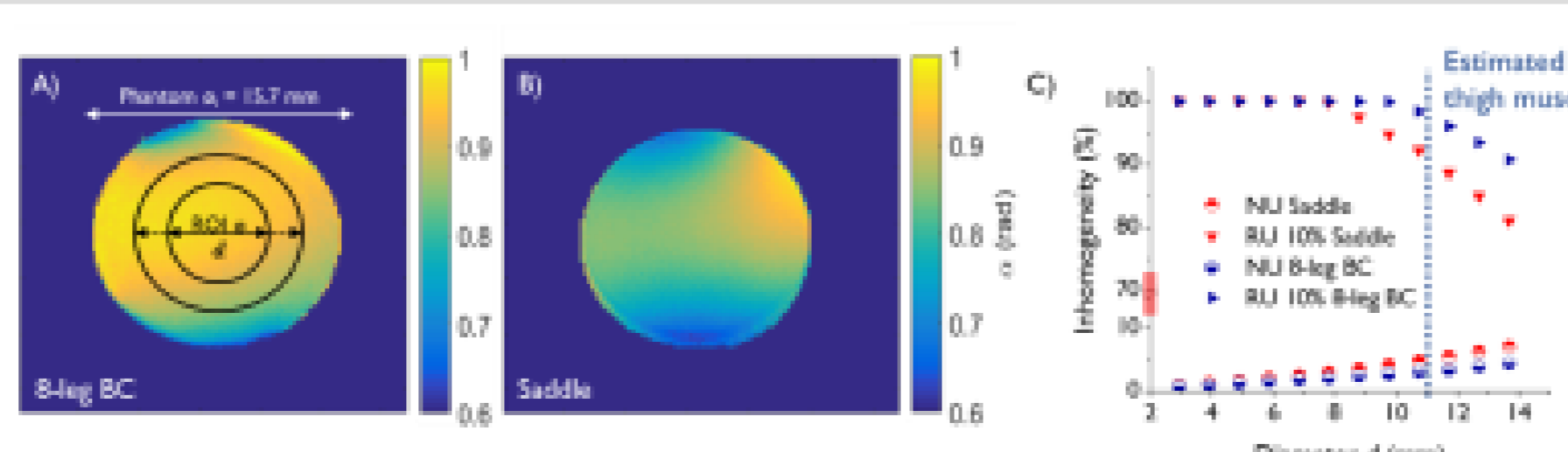


Figure 4: Transmit field homogeneity analysis. Normalized B_1^+ maps of a 10 ml 0.9% saline phantom (inner diameter $\phi_i = 15.7$ mm) generated by **A)** 8-leg BC and **B)** saddle coil. B_1^+ maps were created using the double angle method. Circles represent ROI positioning chosen in C) for statistics. **C)** B_1^+ homogeneity was evaluated via RU and NU as a function of ROI diameter. Acquisition was done with a gradient echo sequence using $TR / TE = 20000 / 4.83$ ms, $\theta = 60^\circ / 120^\circ$, matrix size = 128×128 , $FOV = 25 \times 25$ mm².

	Q_{ul}	Q_l	Q_{ul}/Q_l
Saddle	415	155	2.7
8-leg	84	78	1.1

Table 1: Quality factors for both coils in unloaded (Q_{ul}) and loaded (Q_l) state with a 0.9% saline phantom ($\phi = 15.7$ mm). Q -factors were measured on using the network analyzer by dividing the center frequency by the -7 dB bandwidth.